Natural Convective Mass Transfer MHD Flow of Chemically Reactive Micropolar Fluid past a Vertical Porous Plate

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ABSTRACT

Magnetic field effects on a free convective mass transfer flow of chemically reactive micropolar fluid over a vertical porous plate are investigated in this work. A mathematical model related to the problem is developed from the basis of studying magnetohydrodynamics(MHD). A usual mathematical transformation is applied on the model to obtain a system of non-dimensional equations. Analytical solution of the dimensionless problem is obtained using perturbation technique. The influence of different parameters (Modified Grashof number G_m ,Suction parameter f_w, Magnetic force number M, permeability of porous plate K, Micro-rotational number Δ, Vortex viscosity λ, Spin gradient viscosity number, Schmidt number S_c and Chemical reaction parameter C_r) on velocity, angular velocity and species concentration profiles are presented graphically. Based on these curves the results and conclusion are depicted.

Keywords: MHD, Micropolar fluid, Porous plate, Chemical reaction, Perturbation technique.

1. INTRODUCTION

In recent years, the theory of micropolar fluid has fascinated momentous attention among the researchers. Indeed, the theory of micropolar fluids is introduced by Eringen(1960) can be used to explicate the flow characteristics of polymers, paints, animal blood, lubricants, liquid crystals, hematological suspensions, geomorphological sediments and geomorphological sediments etc. The theory described about local effect increasing due to microstructure and the intrinsic motion of the fluid elements. The model of micropolar fluid represents fluids consisting of rigid, randomly generated particles suspended in a viscous medium where the deformation of the particles is neglected. Kucaba-Pietal (2004), Khedr (2009) investigated colloidal, Muthu (2008) analyzed human and animal blood, Lockwood *et al.* (1987) driven liquid crystal as well as exotic lubricants. Kelson and Desseaux (2001) investigated the effect of surface circumstance on micropolar fluid flow. The unsteady micropolar fluid flow between two parallel porous plates was driven Srinivasacharya *et al.* (2001). Mixed convection micropolar fluid on a porous stretching sheet is investigated by Bhargava (2003). Tripathy *et al* (2016) has investigated hydromagnetic micropolar fluid on a porous stretching sheet. Mohanty *et al.* (2015) investigated heat and mass transfer effect on micropolar fluid on a porous stretching sheet. The mathematical theory of equations of micropolar fluids and applications of these fluids in the theory of lubrication and in the theory of porous media has been studied by Lukaszewicz (1999).

Fluid dynamics is the area of science which is related with the study of motion of fluids. Magnetohydrodynamics (MHD) is the study of the flow of electrically conducting fluids in the presence of a magnetic field. The study of Magnetohydrodynamics has numerous applications such as induction flow meter, cooling of nuclear reactors, structure of stars and planets. Recently, the application of Magnetohydrodynamics in the polymer industry and metallurgy has fascinated the attention of many scholars.

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Nomenclature

u, v	Velocity component	D_m	Diffusion coefficient
g	Acceleration	C_r	Chemical reaction parameter
β^*	concentration expansion coefficient	C_∞	Concentration of uniform flow
B_0	Magnetic field component	Δ	Micro-rotational number
χ	Vortex viscosity	G_m	Modified Grashof number
γ	Spin gradient viscosity	K	Permeability of porous plate
υ	Viscosity	M	Magnetic force number
ρ	Density	$\wedge$	Spin gradient viscosity number
S_c	Schmidt number	λ	Vortex viscosity

Free convection flow takes place frequently in nature, flows of fluid through porous media are of significant interest now days and have fascinated by many researchers owing to their applications in the science and technology.Study of fluid flow in porous medium is based upon the empirically Darcy's law. Such flows are assumed to be useful in diminishing the free convection. Along with the free convection flow, the phenomenon of mass transfer is also very significant in the theories of stellar structure. The study of convective flow with mass transfer along a vertical porous plate is receiving paramount interest of many scholars because of its numerous applications in the field of geophysical and cosmical sciences. Prathapkumer (2010) studied on free convection flow of micropolar and viscous fluids through a vertical duct. Raptis (2011), Samiulhaq*et al.* (2012) and Seth *et al.* (2015) studied on free convective oscillatory flow and mass transfer with ramped temperature on a porous plate. Thereafter, Chamkha (2000), Chaudhary (2007), Samad and Mohebujjaman (2009), Eldabe (2011) and Seth (2015) have paid attention to the study of MHD free convection and mass transfer flows.

The chemical reaction rate counts on the concentration of the species itself. In many chemical engineering systems, there is the chemical reaction between a foreign mass and the fluid. These system occurs in many industrial applications such aspolymer production, manufacturing of ceramics and food processing. Mishra *et al* (2016) has analyzed the effect of chemical reaction on hydromagnetic micropolar fluid flow. Raju *et al.* (2013) investigated an unsteady free convection and chemically reactive MHD flow through infinite vertical porous plate. Bakr (2011) has driven the effect of chemical reaction on micropolar fluid with oscillatory plate. Das *et al.* (1994) analyzed the effective of first order homogeneous chemical reaction of an unsteady micropolar fluid flow. Ibrahim *et al.* (2008),Anand Rao *et al*. (2012), Das (2012) and Raju *et al.* (2013) investigated the effect of chemical reaction on an unsteady MHD free convection fluid through semi-infinite vertical porous plate with heat absorption.Bakr (2011) analyzed the characteristic of a micropolar fluid velocity on oscillatory plate and constant heat source in a rotating frame. Kucaba-Pietal (2004), Khedr*et al*. (2009) investigated the micro inertia effects on the flow of a micropolar fluid past a semi-infinite plate.

Hence our main goal is to investigate a free convective mass transfer steady flow of a chemically reactive micropolar fluid past a semi-infinite porous plate.

2. FORMULATION OF THE PROBLEM

A natural convective mass transfer steady flow of a chemically reactive micropolar fluid through a semi-infinite vertical porous plate is taken into account at the presence of magnetic field. The flow is considered vertically by x-direction and y-direction is represented horizontally. When the flow at rest, the species concentration level $C=C_\infty$ at all point, where C_∞ be the concentration of uniform flow. It is also assumed that a magnetic field B of uniform strength is applied normal to the flow region. The physical configuration and co-ordinate system of the problem is presented in the following Fig.1.

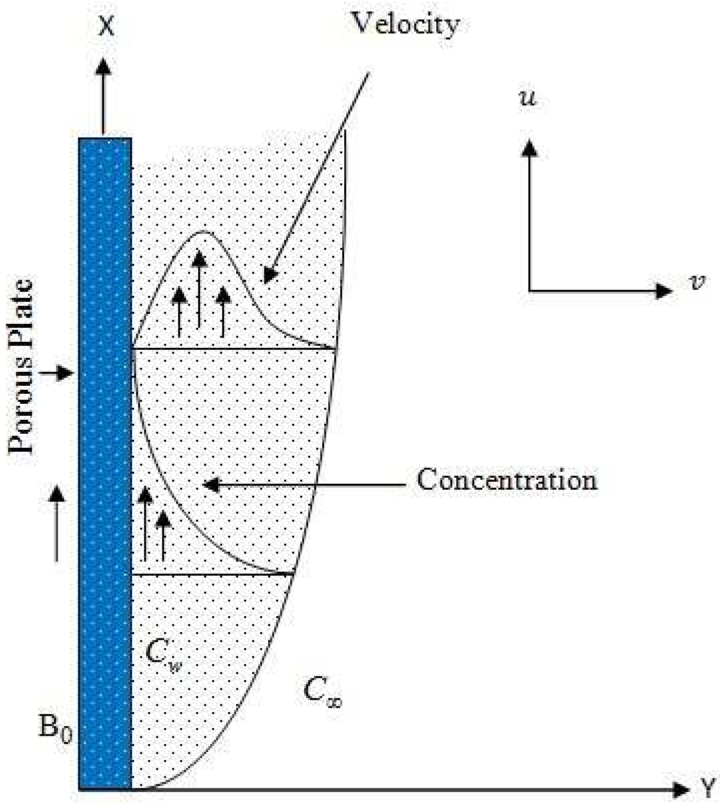


Fig. 1. Physical configuration of the flow

Within the framework of the above stated assumption, the governing equations under the boundary-layer approximations are given by,

Continuity Equation $\dfrac{\partial u}{\partial x}+\dfrac{\partial v}{\partial y}=0$ (1)

Momentum Equation

$$u\frac{\partial u}{\partial x}+v\frac{\partial u}{\partial y}=g\beta^{*}\left(C-C_{\infty}\right)+\left(\upsilon+\frac{\chi}{\rho}\right)\frac{\partial^{2}u}{\partial y^{2}}+\frac{\chi}{\rho}\frac{\partial\Gamma}{\partial y}-\frac{\sigma' uB_{0}^{2}}{\rho}-\frac{\upsilon}{K_{1}}u \tag{2}$$

Angular Momentum Equation

$$u\frac{\partial\Gamma}{\partial x}+v\frac{\partial\Gamma}{\partial y}=\frac{\gamma}{\rho j}\left(\frac{\partial^{2}\Gamma}{\partial y^{2}}\right)-\frac{\chi}{\rho j}\left(2\Gamma+\frac{\partial u}{\partial y}\right) \tag{3}$$

Concentration Equation

$$u\frac{\partial C}{\partial x}+v\frac{\partial C}{\partial y}=D_{m}\frac{\partial^{2}C}{\partial y^{2}}-K_{c}\left(C-C_{\infty}\right) \tag{4}$$

with boundary condition,

$$u=0,\ \Gamma=-s\frac{\partial u}{\partial y},\ C=C_{w}\ \text{ at } y=0$$

$$u=0,\ \Gamma=0,\qquad C=C_{\infty}\ \text{ at } y\to\infty$$

where u is the velocity component, Γ is the velocity acting in z - direction (the rotation of Γ is in the $x-y$ plane), B_0 is the magnetic field component, g is local acceleration due to gravity, χ is the vortex viscosity, γ is the spin gradient viscosity, β^{*} is concentration expansion coefficient.

2.1 Non-dimensional Form

Since our goal is to attain analytical solutions of the problem so we introduce the following dimensionless variables,

$$\eta = y\sqrt{\frac{U_0}{2\upsilon x}}\ ,\psi = \sqrt{2\upsilon U_0 x}\, f(\eta), u = \frac{\partial\psi}{\partial y} \text{ and } v = -\frac{\partial\psi}{\partial x},\ \phi(\eta) = \frac{C - C_\infty}{C_w - C_\infty} \text{ and } \Gamma = \sqrt{\frac{U_0^{\ 3}}{2\upsilon x}}\, g(\eta)$$

The dimensionless equations are,

$$(1+\Delta) f'''(\eta) + \Delta g'(\eta) + f(\eta) f''(\eta) + G_M\varphi(\eta) - (K+M) f'(\eta) = 0 \tag{5}$$

$$\wedge g''(\eta) + f'(\eta) g(\eta) + f(\eta) g'(\eta) - 2\lambda g(\eta) - \lambda f''(\eta) = 0 \tag{6}$$

$$\varphi''(\eta) + S_c f(\eta)\varphi'(\eta) - S_c C_r \varphi(\eta) = 0 \tag{7}$$

The associate boundary conditions

$$f(\eta) = f_w,\ f'(\eta) = 0,\ g(\eta) = -sf''(\eta),\ \varphi(\eta) = 1\ \ \eta = 0$$

$$f'(\eta) = 0,\ g(\eta) = 0,\ \ \varphi(\eta) = 0\ \ \eta \to \infty$$

Where, micro-rotational number $\Delta = \frac{\chi}{\rho\upsilon}$, modified Grashof number, $G_m = \frac{g\beta^*(C_w - C_\infty)2x}{U_0^{\ 2}}$ permeability of porous plate, $K = \frac{2\upsilon x}{K_1 U_0}$, magnetic force number, $M = \frac{\sigma B_0^{\ 2} 2x}{U_0\rho}$, Spin Gradient number, $\wedge = \frac{\gamma}{\upsilon\rho j}$, Vortex viscosity, $\lambda = \frac{2x\chi}{\rho j U_0}$, Schmidt number, $S_c = \frac{\upsilon}{D_m}$, chemical reaction parameter, $C_r = K_c\frac{2x}{U_0}$.

3 ANALITICAL SOLUTION

Since the solution is sought for the large suction further transformation can be made **Arifuzzaman (2015)** as,

$$\xi = \eta f_w \tag{7}$$

$$f(\eta) = f_w F(\xi) \tag{8}$$

$$\phi(\eta) = f_w^{\ 2} G(\xi) \tag{9}$$

$$g(\eta) = f_w^{\ 3} H(\xi) \tag{10}$$

Now the model with small quantity,

$$(1+\Delta)F''' + \Delta H' + FF'' + \varepsilon G_M G - (K+M)\varepsilon F' = 0 \tag{11}$$

$$\wedge H'' + F'H + FH' - 2\lambda\varepsilon H - \lambda\varepsilon F'' = 0 \tag{12}$$

$$G'' + S_c G'F - \varepsilon S_c C_r G = 0 \tag{13}$$

The associate boundary conditions

$$F(\xi) = 1,\quad F'(\xi) = 0,\ G(\xi) = \varepsilon,\ H(\xi) = -sF''(\xi) \text{ at } \xi = 0$$

$$F'(\xi) = 0,\ G(\xi) = 0,\quad H(\xi) = 0 \text{ at } \xi \to \infty$$

Now for the large suction $(f_w > 1)$, ε will be very small. Therefore following **Bestman** (1990), F, G and H can be expended in terms of the small perturbation quantity ε,

$$F(\xi) = 1 + \varepsilon F_1(\xi) + \varepsilon^2 F_2(\xi) + \varepsilon^3 F_3(\xi) + \ldots\ldots \tag{14}$$

$$G(\xi) = \varepsilon G_1(\xi) + \varepsilon^2 G_2(\xi) + \varepsilon^3 G_3(\xi) + \ldots\ldots \tag{15}$$

$$H(\xi) = \varepsilon H_1(\xi) + \varepsilon^2 H_2(\xi) + \varepsilon^3 H_3(\xi) + \ldots\ldots \tag{16}$$

The dimensionless equations (5) - (7) transform into the following first order, second order and third order equations with their associated boundary conditions,

$$\left.\begin{aligned}&\left(1+\Delta\right)F_1''' + \Delta H_1' + F_1'' = 0\\ &\wedge H_1'' + H_1' = 0\\ &G_1'' + S_c G_1' = 0\end{aligned}\right\} \tag{17}$$

$F_1 = 0,\ \ G_1 = 1,\ \ H_1 = 0$ at $\xi = 0$

$F_1' = 0,\ \ G_1 = 0,\ H_1 = 0$ at $\xi \to \infty$

$$\left.\begin{aligned}&\left(1+\Delta\right)F_2''' + \Delta H_2'\left(\xi\right) + F_1 F_1'' + F_2'' + G_M G_1\left(\xi\right) - \left(K+M\right)F_1' = 0\\ &\wedge H_2'' + F_1' H_1 + F_1 H_1' + H_2' - 2\lambda H_1 - \lambda F_1'' = 0\\ &G_2'' + S_c G_2' + S_c G_1' F_1 - S_c C_r G_1 = 0\end{aligned}\right\} \tag{18}$$

$F_2 = 0,\qquad G_2 = 0,\ \ H_2 = 0$ at $\xi = 0$

$F_2' = -\dfrac{1}{\varepsilon^2},\ G_2 = 0,\ \ H_2 = 0$ at $\xi \to \infty$

$$\left.\begin{aligned}&\left(1+\Delta\right)F_3''' + \Delta H_3' + F_2 F_1'' + F_1 F_2'' + F_3'' + G_M G_2 - \left(K+M\right)F_2' = 0\\ &\wedge H_3'' + F_2' H_1 + F_1' H_2 + F_2 H_1' + F_1 H_2' + H_3' - 2\lambda H_2 - \lambda F_2'' = 0\\ &G_3'' + S_c G_3' + S_c G_2' F_1 + S_c G_1' F_2 - S_c C_r G_2 = 0\end{aligned}\right\} \tag{19}$$

$F_3 = 0,\ \ G_3 = 0,\ \ H_3 = 0$ at $\xi = 0$

$F_3' = \dfrac{1}{\varepsilon^3},\ G_3 = 0,\ \ H_3 = 0$ at $\xi \to \infty$

Now the solution of first order, second order and third order equations are given following,

$F_1 = 0,\ \ G_1 = e^{-S_c\xi},\ \ H_1 = 0$

$F_2 = \dfrac{1}{A_2\varepsilon^2} - \dfrac{\xi}{\varepsilon^2} - \dfrac{1}{A_2\varepsilon^2}e^{-A_2\xi} + A_3 e^{-S_c\xi}$, $H_2 = 0$ and $G_2 = -C_r e^{-S_c\xi}$

$F_3 = -\dfrac{1}{A_2\varepsilon^3} + \dfrac{\xi}{\varepsilon^3} + \dfrac{1}{A_2\varepsilon^3}e^{-A_2\xi} + A_{15}e^{-A_2\xi} + A_{16}e^{-S_c\xi} + A_{17}e^{-S_c\xi} + A_{18}e^{-A_2\xi} - A_{19}e^{-S_c\xi}\ \ H_3 = A_8 e^{-A_2\xi} + A_9 e^{-S_c\xi}$

$G_3 = A_4 e^{-S_c\xi} - A_5 e^{-S_c\xi} + A_6\xi e^{-S_c\xi} - A_7 e^{-\left(S_c + A_2\right)\xi} + A_{20}e^{-2S_c\xi}$

Substituting the values of *F*, *H* and *G*, We get

The velocity equation,

$$f = f_w F\left(\xi\right)$$

$$\Rightarrow f = f_w + \left(\varepsilon^2 f_w A_3 + \varepsilon^3 f_w A_{16} + \varepsilon^3 f_w A_{17}e^{-S_c\xi} - \varepsilon^3 f_w A_{19}\right)e^{-S_c\xi} + \left(\varepsilon^3 f_w A_{15} + \varepsilon^3 f_w A_{18}\right)e^{-A_2\xi}$$

$$\Rightarrow f' = \left(-S_c\varepsilon^2 f_w A_3 - S_c\varepsilon^3 f_w A_{16} - S_c\varepsilon^3 f_w A_{17}e^{-S_c\xi} + S_c\varepsilon^3 f_w A_{19}\right)e^{-S_c\xi}$$

$$-\left(A_2\varepsilon^3 f_w A_{15} + A_2\varepsilon^3 f_w A_{18}\right)e^{-A_2\xi}$$

The angular velocity equation,

$$g\left(\eta\right) = f_w^3 H\left(\xi\right)$$

$$\Rightarrow g(\eta) = \varepsilon^3 f_w^3 A_8 e^{-A_2\xi} + \varepsilon^3 f_w^3 A_9 e^{-S_c\xi}$$

The concentration equation,

$$\phi(\eta) = f_w^{\ 2} G(\xi)$$

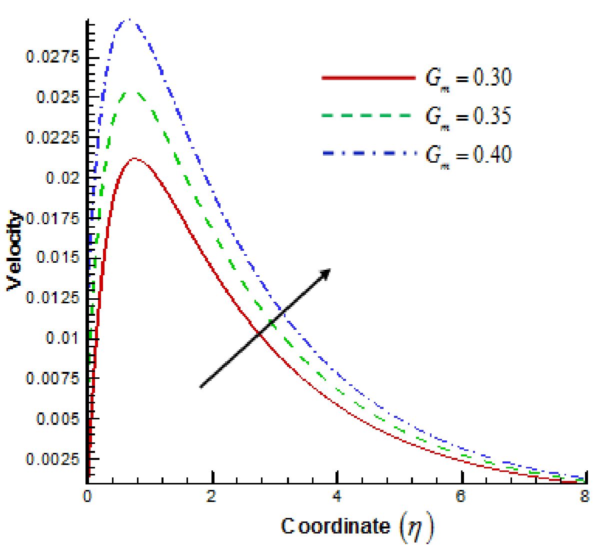
$$\Rightarrow \varphi(\eta) = f_w^{\ 2} \varepsilon e^{-S_c\xi} - \varepsilon^2 f_w^{\ 2} C_r e^{-S_c\xi} + \varepsilon^3 f_w^{\ 2} A_4 e^{-S_c\xi} - \varepsilon^3 f_w^{\ 2} A_5 e^{-S_c\xi} + \varepsilon^3 f_w^{\ 2} A_6 \xi e^{-S_c\xi}$$
$$-\varepsilon^3 f_w^{\ 2} A_7 e^{-(S_c+A_2)\xi}$$

4. RESULT AND DISCUSSION

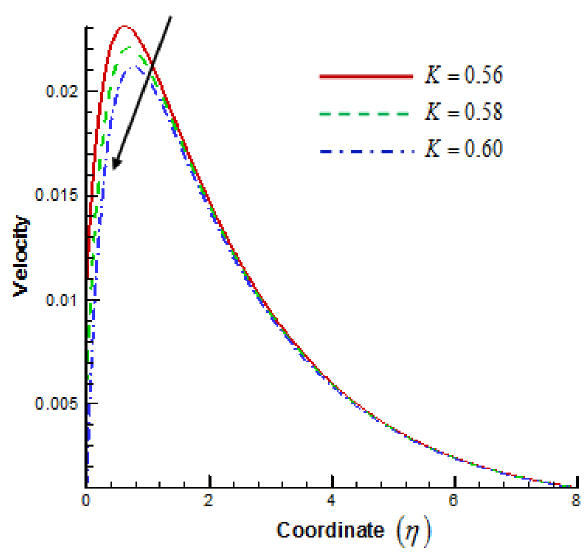
For the purpose of the applicability of the present mathematical model, the analytical solution are driven using the perturbation method and the discussion is made for various values of parameters just as Modified grashof number G_m ,Suction parameter f_w, Magnetic force number M, permeability of porous plate K, Micro-rotational number Δ, Vortex viscosity λ, Spin gradient viscosity number $(\wedge)$, Schmidt number S_c and Chemical reaction parameter C_r. The fluid velocity, angular velocity and concentration versus the non-dimensional coordinate variable η are displayed in Figures.

The increase values of magnetic parameter create a drag force known as Lorentz force. The velocity profiles are illustrated in Fig. 2. As it is observed, the velocity profiles curve climb up at the increase of magnetic force number. Afterwards, the suction parameter (f_w) stabilizes the boundary layer growth. So the velocities profiles curve decline with go up suction parameters. Schmidt number decreases the molecular diffusivity. Modified Grashof number signifies the effect of buoyancy force to the viscous hydrodynamic force. So, velocity curves increase with modified Grashof number (G_m). The velocity profiles go down for permeability of porous plate (K). Then with increase of vortex viscosity, the velocity profiles plunge. Fig. 3 revels the angular velocity profiles. Firstly, angular velocity profiles decline with rise of suction parameter (f_w). After that, it increases with the increase of modified Grashof number (G_m). But in Fig. 3(c) angular velocity profiles show fluctuation for Schmidt number (S_c). Then at the upsurge of micro-rotational number the angle velocity increase. At the end of the list of figure, angular velocity decline due to soar of spin gradient viscosity number and vortex viscosity.

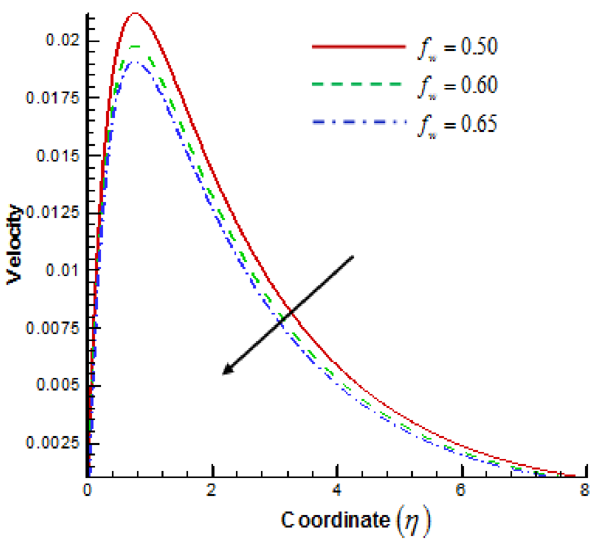
Fig. 4 describes the concentration profiles. As it is noticed, concentration boundary layer is lowered down as chemical reaction parameter (C_r) climb up. Concentration curves are also declined as Schmidt number (S_c), suction parameter (f_w), micro-rotational parameter Δ increase.



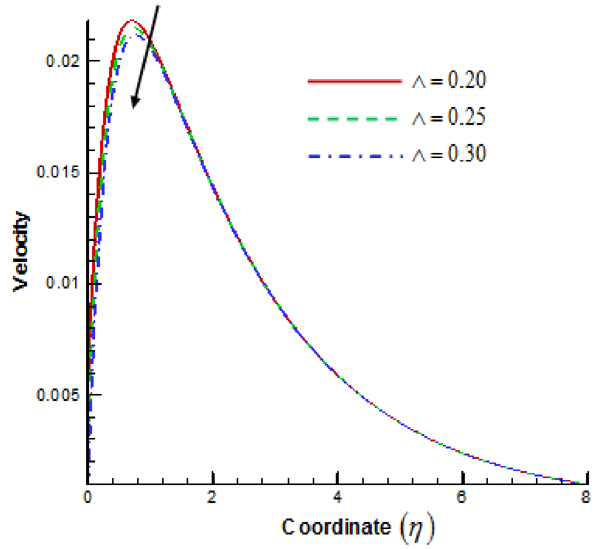
(a)



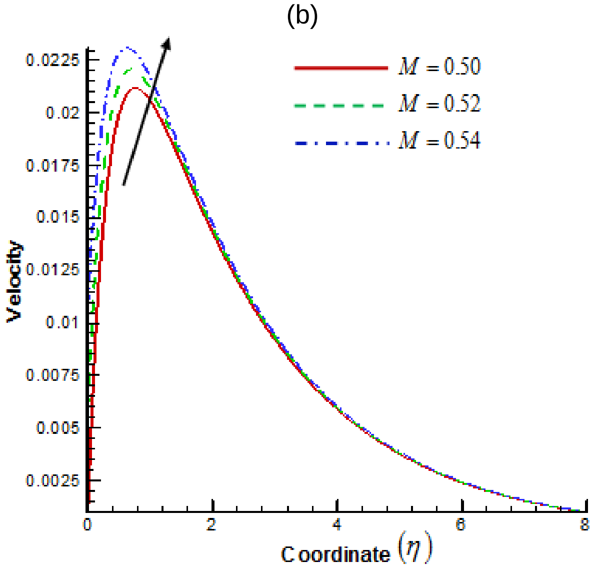
(d)



(b)

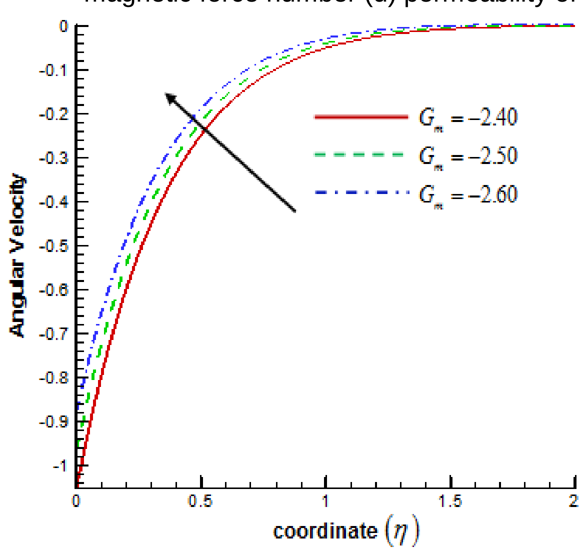


(e)

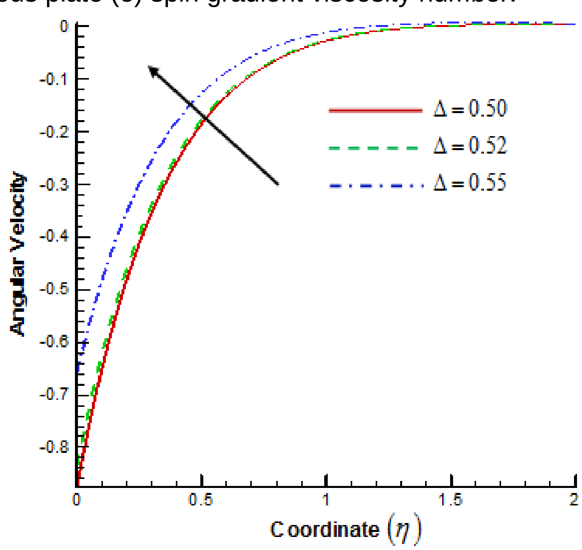


(c)

Fig. 2. Velocity profiles for different values of (a) modified Grashof number (b) suction parameter (c) magnetic force number (d) permeability of porous plate (e) spin gradient viscosity number.



(a)



(d)

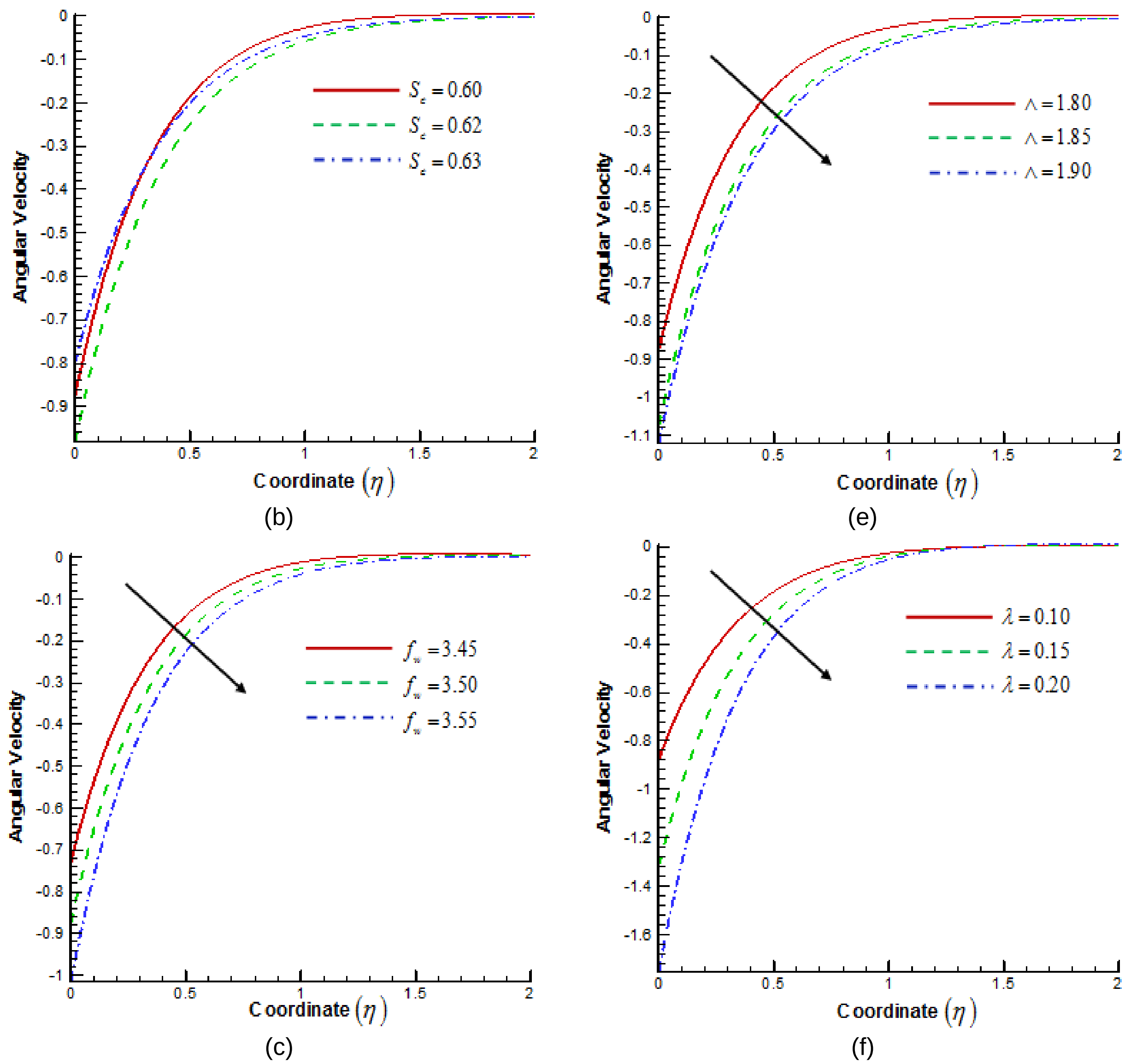
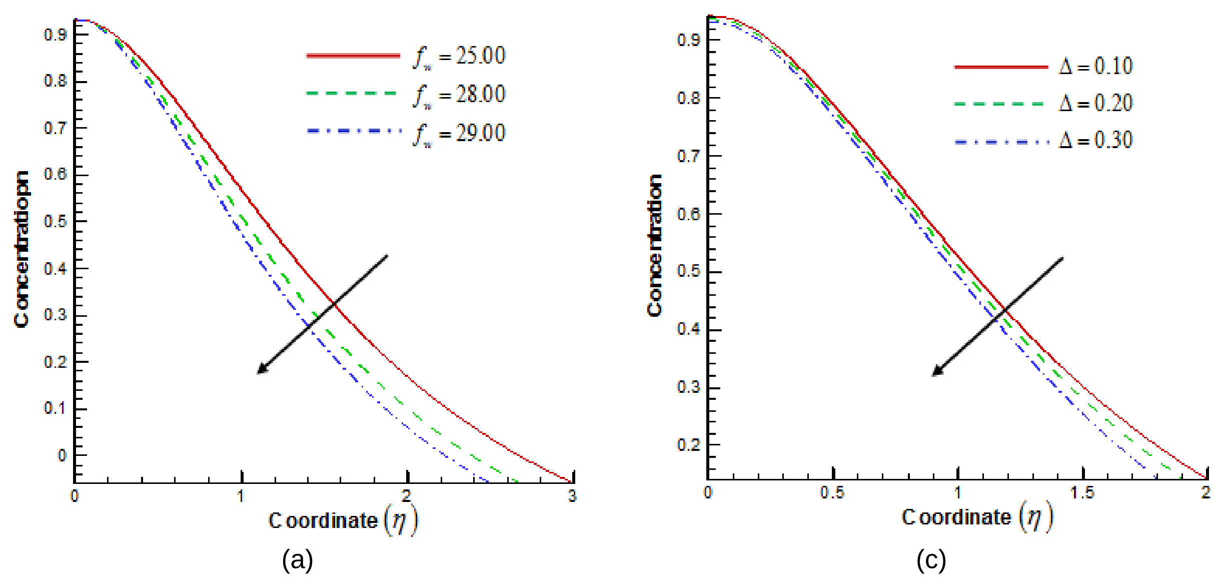
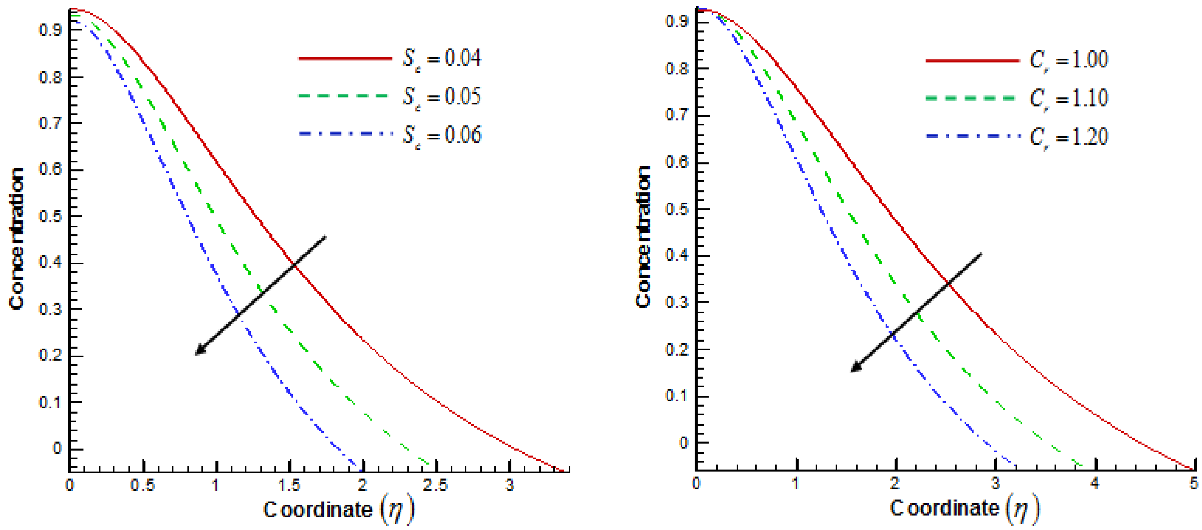


Fig. 3. Angular velocity profiles for different values of (a) modified Grashof number (b) Schmidt number (c) suction parameter (d) micro-rotational number (e) spin gradient viscosity number (f) vortex viscosity.





(b) (d)

Fig. 4. Concentration profiles for different values of (a) suction parameter (d) Schmidt number (c) micro-rotational number (d) chemical reaction parameter.

4. CONCLUSIONS

Some of the important findings of the present work obtained from the graphical representation of the results are listed below:

1. The fluid velocity and angular velocity profile decreases with the increase of Modified Grashof number.
2. The velocity and angular velocity profile decreases with the increase of Suction parameter and also the concentration profile decreases with the increase of Suction parameter.
3. The velocity profile decreases and angular velocity profiles decreases with the increase of Schmidt number and also the concentration profile decreases with the increase of Schmidt number.
4. The velocity profile increase with Magnetic force number.
5. The velocity profiles decreases with the increase of Permeability of porous plate.
6. The concentration profile decreases with the increase of Chemical reaction parameter.
7. The angular velocity profile increases with the increase of Micro-rotational number and the concentration profile decreases with the increase of Micro-rotational number.
8. The angular velocity profile decreases with the increase of Spin gradient viscosity number.
9. The angular velocity profile decreases with the increase of Vortex viscosity.

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APPENDIX

$$A_1 = \frac{1}{\wedge}, \quad A_2 = \frac{1}{1+\Delta}, \quad A_3 = \frac{-A_2 G_M}{\left(-S_c{}^3 + A_2 S_c{}^2\right)}, \quad A_5 = \frac{S_c}{A_2 \varepsilon^2}, \quad A_6 = \frac{S_c}{\varepsilon^2} \quad A_7 = \frac{S_c^2}{A_2 \varepsilon^2 \left(A_2{}^2 + S_c A_2\right)},$$

$$A_{20} = \frac{A_3}{2}, \quad A_8 = \frac{-\frac{A_2 \lambda}{\varepsilon^2}}{\left(A_2^2 - A_1 A_2\right)}, \quad A_9 = \frac{A_3 \lambda S_c^2}{S_c^2 - A_1 S_c}, \quad A_{10} = \Delta A_2{}^2 A_8 \quad A_{11} = \Delta A_2 A_9 S_c, \quad A_{12} = G_m C_r A_2,$$

$$A_{13} = \frac{A_2\left(K+M\right)}{\varepsilon^2}, \quad A_{14} = A_2 S_c A_3\left(K+M\right), \quad A_{15} = \frac{A_{10}}{A_2{}^2}, \quad A_{16} = \frac{A_{11}}{\left(-S_c^3 + S_c^2 A_2\right)}, \quad A_{17} = \frac{A_{12}}{\left(-S_c^3 + S_c^2 A_2\right)},$$

$$A_{18} = \frac{A_{13}}{A_2{}^2}, \quad A_{19} = \frac{A_{14}}{\left(-S_c^3 + S_c^2 A_2\right)}$$