

A New Quantum Paradox

Abstract: A gauge transformation of a simple electromagnetic system is analyzed. The Hamiltonian which is derived from the Dirac Lagrangian density is used for determining the state of an electron. The fact that this Hamiltonian is free of time differential operators plays a key role in the analysis. It is proved that an application of a specific gauge transformation yields inconsistent results. These results call for a further analysis of the role of gauge transformations in the theoretical structure of electrodynamics.

Keywords: Lagrangian density, Dirac equation, Hamiltonian, Gauge transformations, Paradox

1. Introduction

This work discusses a paradox that is obtained from an application of a gauge transformation to a simple electromagnetic system. Electrodynamics is a widely studied sector of theoretical physics and it is relevant to many physical disciplines, ranging from solid state physics to astrophysics. Therefore, the entire physical community is expected to belong to the readership of this work.

A paradox is regarded as a useful tool for finding out new properties and interrelations between elements of a theory. A physical paradox describes a hypothetical device and relevant physical laws are assumed to determine the behavior of the system. The outcome of a paradox is an apparent contradiction. Such a contradiction provides a motivation for a further investigation of the relevant physical laws. This kind of investigation generally contributes to a better understanding of these laws. The following lines briefly describe two well known paradoxes which are used here as an illustration of this matter.

In the 1930s, Einstein, Podolsky and Rosen (EPR) described a quantum paradox of an action at a distance [1]. They used a principle which they called *physical reality* and regarded the result as an indication that quantum mechanics is an incomplete theory. For this reason, EPR put forward the need for finding hidden parameters that will promote quantum mechanics to the status of a complete theory. Later Bohm and Aharonov [2, 3] and Bell [4] have added elements that were used in an experimental test of the EPR idea. Experimental results support the idea that there is a kind of quantum information that propagates instantaneously (see [5] and references therein). Thus, the apparent EPR paradox has provided a motivation for acquiring new information on how physical processes work.

In the 1960s Shockley and James presented a paradox where a stationary system of a charge and a magnet has an electromagnetic nonzero linear momentum [6]. Soon

after the publication of this paradox, Coleman and Van Vleck provided a general proof showing that the system's total linear momentum must be balanced [7]. Later Comay has shown that an explicit mechanical linear momentum exists in the system. In particular, if a nonvanishing pressure gradient exists along a closed loop of current then effects related to the energy-momentum tensor yield a nonzero mechanical linear momentum [8]. This mechanical momentum balances the electromagnetic linear momentum and also supports the validity of Coleman and Van Vleck general analysis. Thus, the Shockley and James paradox has ended up with a better understanding of elements of classical physics.

This paper discusses gauge transformations in the quantum domain. In classical physics, electrodynamic equations of motion - namely, Maxwell equations and the Lorentz force - are independent of the 4-potentials. Therefore, classical electrodynamics is invariant under a gauge transformation. On the other hand, quantum theories depend explicitly on the 4-potentials. The analysis abides by physical laws and proves that the Dirac Hamiltonian is not invariant under a specific gauge transformation. This outcome demonstrates the need for a further analysis of the role of gauge in theoretical physics.

The paradox of this work is described in the second section. The third section contains the conclusions. Expressions are written in units where $\hbar = c = 1$. The relativistic metric is diagonal and its entries are (1,-1,-1,-1). Greek indices run from 0 to 3.

2. The Paradox

The paradox described below arises from an examination of a specific gauge transformation that pertains to the state of an electron which obeys the Dirac equation.

To this end, let us examine the Lagrangian density of a Dirac electron [9, see p. 78]

$$\mathcal{L}_D = \bar{\psi}[\gamma^\mu(i\partial_\mu - eA_\mu) - m]\psi, \quad (1)$$

where $A^\mu = (V, \mathbf{A})$ denote the components of the electromagnetic 4-potential [10, see p. 10] or [11, see p. 48]. Here one sees that in this equation, like in any other quantum equation, the charge interacts with the 4-potential.

It is well known that the Lagrangian density of (1) is invariant under the gauge transformation $\Lambda(x)$ which is an arbitrary function of the space-time coordinates (denoted by x) [9, see p. 78] and [12, see p. 345]

$$A_\mu(x) \rightarrow A_\mu(x) + \Lambda(x)_{,\mu}; \quad \psi(x) \rightarrow \exp(-ie\Lambda(x))\psi(x). \quad (2)$$

Here e is the electronic charge, which is a dimensionless Lorentz scalar in the units where $\hbar = c = 1$. Indeed, substituting (2) into (1), one finds that the contribution of the gauge 4-potentials $\Lambda(x)_{,\mu}$ is canceled out by the additional terms obtained from the partial differentiation of $\exp(-ie\Lambda(x))\psi(x)$.

Let us turn to the paradox and examine a motionless electron located at the vicinity of point P in a field-free space. The Dirac Hamiltonian is used for finding the time evolution of this electron. (A quantum expression for the Hamiltonian is also required by the Bohr correspondence principle. Here the classical limit of quantum theories should agree with classical physics. Evidently, in classical physics energy is a well defined quantity. Therefore, one requires that quantum theories should have a self-consistent expression for energy.) This Hamiltonian can be derived from the Lagrangian density of (1) in the following steps.

The Hamiltonian density \mathcal{H} is derived from the Lagrangian density by the well known Legendre transformation

$$\mathcal{H} = \sum_i \frac{\partial \mathcal{L}}{\partial \dot{\psi}_i} \dot{\psi}_i - \mathcal{L}, \quad (3)$$

where the index i runs on all functions. In the specific case of a Dirac particle one obtains from (1) and (3)

$$\mathcal{H}_D = \psi^\dagger [\boldsymbol{\alpha} \cdot (-i\nabla - e\mathbf{A}) + \beta m + eV] \psi, \quad (4)$$

which is written here in the standard notation [10, see p. 11]. The density of a Dirac particle is $\psi^\dagger \psi$ [10, see p. 9]. Thus, removing the density from (4), one obtains the operator form of the Dirac Hamiltonian

$$H_D = [\boldsymbol{\alpha} \cdot (-i\nabla - e\mathbf{A}) + \beta m + eV]. \quad (5)$$

This Hamiltonian stands on the right hand side of the Dirac equation [10, see p. 11]

$$i \frac{\partial \psi}{\partial t} = H_D \psi = [\boldsymbol{\alpha} \cdot (-i\nabla - e\mathbf{A}) + \beta m + eV] \psi. \quad (6)$$

As is well known, the Dirac Hamiltonian (5) does not contain a time differential operator.

The Dirac equation (6) is used for finding the time evolution of an electron at the vicinity of point P . Here the field-free 4-potential is

$$A^\mu = 0. \quad (7)$$

Hence, the Dirac equation for a free electron

$$i \frac{\partial \psi}{\partial t} = [\boldsymbol{\alpha} \cdot (-i\nabla) + \beta m] \psi \quad (8)$$

determines the electronic state.

Let us examine how this system is affected by the following gauge function

$$\Lambda(x) = et/r. \quad (9)$$

Here e is the absolute value of the electronic charge, t is the time and r is the distance from the origin of the spatial coordinates. Certainly, the gauge function of (9) is a

legitimate gauge expression because it depends on space-time coordinates. This gauge transformation casts the null 4-potential of (7) into the following expression

$$A'_\mu = \frac{\partial(et/r)}{\partial x^\mu} = (e/r, -et\mathbf{r}/r^3). \quad (10)$$

The gauge transformation of (2) also transforms the Dirac wave function. Introducing the specific gauge function of (9), one finds that the transformed wave function (2) is

$$\psi'(x) = \exp(ie^2t/r)\psi(x). \quad (11)$$

And indeed, substituting (10) and (11) into the Dirac Lagrangian density (1), one finds that this Lagrangian density is invariant under the transformation of the gauge function (9).

On the other hand, the new function $\psi'(x)$ of (11) must satisfy the Dirac equation (6), where the gauge terms of (10) are used in the expression for the 4-potential. Here one obtains

$$\begin{aligned} i\frac{\partial\psi'}{\partial t} = H_D\psi'(x) &= [\boldsymbol{\alpha} \cdot (-i\nabla - e\mathbf{A}) + \beta m + eV]\psi'(x) \\ &= \exp(ie^2t/r)[\boldsymbol{\alpha} \cdot (-i\nabla) + \beta m - e^2/r]\psi(x) \end{aligned} \quad (12)$$

It turns out that similarly to the case of the Dirac Lagrangian density (1), the contribution of the 3-vector part of the gauge (10) is eliminated from the Dirac Hamiltonian of (12). On the other hand, the 0-component of that gauge remains as is. This outcome stems from the fact that the Dirac Hamiltonian (5) contains spatial differential operators but is free of a time differential operator.

Now, let us define the point P of the electron so that its distance from the origin of the spatial coordinates is about the Bohr radius. In this case the gauge transformation (9) produces (12), which is the equation of a bound electron of the hydrogen atom (multiplied by a phase factor) [10, see p. 52]. This outcome is inconsistent with the free electron state of the null potential (7).

There is another problem with the gauge transformed function ψ' of (11). Thus, an ordinary wave function of a motionless particle in a well-defined energy state takes the form

$$\psi(x) = \exp(-iEt)\chi(x, y, z). \quad (13)$$

Here the time dependence of $\psi(x)$ appears only in the phase where the energy E is a constant and $\chi(x, y, z)$ is a spatially dependent energy eigenfunction. On the other hand, the phase factor of the gauge transformed function ψ' of (11) also depends on the radial coordinate r . Thus, one obtains for the motionless free electron

$$\begin{aligned} i\frac{\partial\psi'}{\partial t} &= i\frac{\partial\exp(ie^2t/r)}{\partial t}\psi + i\exp(ie^2t/r)\frac{\partial\psi}{\partial t} \\ &= (-e^2/r + m)\psi', \end{aligned} \quad (14)$$

where m is the electronic mass. Therefore, the coordinate-dependent quantity $-e^2/r$ of (14) proves that the gauge-transformed wave function ψ' is *not* an energy eigenfunction. This is a contradiction because an electron in a free space has a well defined energy and in the case of a motionless electron $E = m$ [10, see p. 28].

3. Conclusions

The paradox described herein is related to two expressions of the 4-potential (7) and (10) which differ by a gauge transformation. It turns out that contrary to a general expectation, two physically different results are obtained. The state of a free electron which is derived from (8) is inconsistent with that of the solution of (12), where the electron is bound to the hydrogen atom. An examination of the derivation of the final form of (12) indicates that the problem stems from the fact that unlike the Dirac Lagrangian density of (1), the Dirac Hamiltonian density of (4) and its associated Hamiltonian of (5) are independent of a time differential operator. The

result of this work provides a motivation for a further investigation of the role of gauge in electrodynamics. Such an investigation is expected to end up with a deeper understanding of the role of gauge in electrodynamics.

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