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#### ELECTRON DIFFRACTION AND ELECTRON MICROSCOPYSTUDY OF FILMS CuGaS<sub>2</sub>.

7 8

#### 9 Abstract

10 In present work, the results of electron diffraction investigations of structures of 11 amorphous thin films of CuGaS<sub>2</sub> have given and function of radial distribution of atoms 12 (FRDA) have been calculated. Appropriate coordination number n=4,1 we obtained 13 from calculating the area under the first peak, also indicates tetrahedral surrounded by 14 atoms of copper and gallium. During the deposition of this ternary compound on a substrate 15 with T = 423-433 K is formed by a mixture of polycrystalline single crystal. With the increase 16 of temprature the intensivity of polycrystallines decreases and point reflections according to the 17 monocrystal increases. Further increase of the substrate temperature to 453 K LiF leads to the 18 formation of a perfect single crystal.

19 Superstructure phase CuGaS<sub>2</sub> is oriented (100) plane parallel to the faces LiF. During 20 epitaxial growth on LiF CuGaS<sub>2</sub> one unit cell superstructure is mated with four cells of the 21 substrate. Between periods of lattices of the initial phase and superstructure there 22 are simple relations common with:  $a \approx 3a_0$ ;  $c \approx 2c_0$ .

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Keywords. Diffraction, phase, atoms, structure, superstructure, amorphous, compositions

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#### 27 **1.Introduction.**

A number of works [1-3] were devoted to X- ray studies of crystalline structures of the compounds of group  $A^1$ - $B^3$ - $C^6$ . However in none of the known works patterns of short range order structure of amorphous compositions CuGaS<sub>2</sub> were determined.

The reason for this can be found either by difficulties in establishing
conditions for amorphous films of these compounds, or trends in amorphous films to the
more dense packing.

35 Amorphous thin films CuGaS<sub>2</sub>,

of thickness 25 nm were obtained by evaporationalloys  $CuGaS_2$  in the vacuum of  $10^{-1}$ 

 $^{4}$  PA on the substrate *NaCl*, *KCl* 

and LiF located at room temperature. The rate of deposition of films for all cases was the

 $11.5 \text{ nm/sec.Amorphous phase CuGaS}_2 \text{ is formed until } T_s = 383 \text{ K}, \text{ crystallization},$ 

40 which canlead to the formation of polycrystalline with periods of a tetragonal lattice,

- 41 military data [5].
- 42 Amorphous filmsformed with values  $S=4\pi \sin\theta/\lambda=24.10$ ; 29.50; 53.70; 83,70 nm<sup>-</sup>

43 <sup>1</sup> (Fig.1)after heat treatment at  $T_s$ =380 K are crystallized in the structure of chalcopyrite

tetragonal lattice CuGaS<sub>2</sub> with periods a=0,535; c=1.047 nm, CBC I $\overline{4}$  2d [5].

- 45 2. Experiment
- 46 Parameters of
- 47 short range order distances of coordination spheres and interatomic distances, coordinatio
- 48 nnumbers (CN) may be determined by the functions of atom radial distribution (FARD) pr
- 49 epared according to

50 
$$4\pi r^2 \rho(r) = 4\pi r^2 \rho_0 + \frac{2r}{\pi} \int_0^\infty Si(S) \sin(sr) dS$$
 (1)[4],

51 the retraining Fourier intensity of the coherent scattering of electrons. Reliable, (FARD) c if 52 an only be obtained when integrating from 0 to  $\infty$  or before S<sub>2</sub>, interference functions do not feel out of oscillate that occurs in strongly disordered 53 systems. 54 The intensity of scattering can be determined experimentally with sufficient accuracy only on some interval of values  $S=4\pi \sin\theta/\lambda$ , so practically the integration 55 in (1)is over a finite interval from  $S_1$  to  $S_2$ . 56

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#### 58 **3. Results and discussion**

Intensity curve electron scattering from amorphous films CuGaS<sub>2</sub> were obtained by us on electronography brand EMR-102 in the form of graphs of dependences of the intensity of scattering angles, i.e. from S= $4\pi \sin\theta/\lambda$  (Fig.1).

62 Function i(S)

graphically depicted in (Figure 2.), were used to calculate (FARD) for CuGaS<sub>2</sub> (fig. 3.) acc ording to formula (1). The calculation was carried out on the programmer"RADIADIS" on the computer IBM. Intervals of variables accounted for  $\Delta r = 0,01 \text{ nm}^{-1}$ ,  $\Delta S = 0,01 \text{ nm}^{-1}$ .

(FARD) CuGaS<sub>2</sub> (fig. 3.) contains four asymmetric highs one of which is isolated and a
group of false highs, emanating from the larger values of S.

The area under the respective highs, manifesting themselves in  $r_1 = 0,234$ ;  $r_2 = 0,244$ ;  $r_3 = 0,282$ ;  $r_4 = 0,411$  nm. are equal  $\Delta_1 = 16,0$ ;  $\Delta_2 = 20,3$ ;  $\Delta_3 = 55,5$  and  $\Delta_4 = 80,0$ nm respectively.

Distance  $r_1 = 0,234$  nm revealed on (FARD) CuGaS<sub>2</sub>, is the distance between atoms Cu-S, as tetrahedral covalent radii are equal to  $r_{Cu} = 0,135$  and  $r_s = 0,104$ nm. Appropriate coordination number n = 4,1 we obtained from calculating the area under the first peak, also indicates tetrahedral surrounded by atoms of copper and gallium.

Radius of the second coordination sphere equal to  $r_2 = 0,244$  nm 77 which we were 78 able to interpret as the distance Ga-S. The ionic radii of gallium and sulfur atoms  $r_{Ga} = 0,127$  and  $r_{S} = 0,104$  nm, for which coordination number is six. 79 constitute Meaning coordination number equal to  $n_2 = 6.3$  received by us from calculating 80 81 the area of the second peak on (FARD) CuGaS<sub>2</sub> also indicates the octahedral environment of gallium atoms by sulfur atoms. 82

to  $r_3 = 0,282 \text{ nm}$ The bond length between the atoms equal 83 corresponds to the 84 distance between the same atoms Ga-Ga. Fourth maximum detected at a distance nm can be referred to the distance between the negatively charged divalent  $atoms(S^{2-})$ . 85 on (FARD) CuGaS<sub>2</sub> arise false details and It should be noted that too. 86 they may arise due to errors in the experimental intensity curve or cliff: they are mainly manifes 87 ted in the large values of S and belong to groups with slightly blurred large 88 89 coordination number equal to 8, 10 and 12.

90 During the deposition of this ternary compound on the substrate with  $T_s = 423$ -

433 K a mixture of polycrystalline with single crystal is formed (Fig. 4.).

92 On the electron from the mixture appears more additional weaker reflexes appear. As a 93 result of increase of temperature intensity polycrystalline lines are reduced but the 94 intensity of point reflections corresponding to single crystal grow.

95	Further increase of the substrate temperature to 453 K LiF leads to the formation of
96	a perfect single crystal. On electron diffraction from single crystals
97	(Fig. 5), discontinued at right angles strong point reflexes, forming a square grid
98	displayed on the basis of hk0 reflexes known lattice CuGaS <sub>2</sub> . Indexing of all
99	reflexes, including additional low in intensity lines is achieved with parameter
100	a = 1.605  nm.
101	Period "c", established by electron diffraction, shot at an angle $\varphi = 35^{\circ}$ was found to be
102	2.102 nm.
103	Between periods of lattices of the initial phase and superstructure there
104	are simple relations common with: $a \approx 3a_0$ ; $c \approx 2c_0$ . In the electron derived from film
105	on substrates formulated withhigher temperatures ( $T_s$ =473 K), dynamic effects
106	appear (Figure 6.). The microstructure of single –
107	crystal layers $CuGaS_2$ from which there is a dynamic high energy electron scattering is sho
108	wn in Fig. 7 (X 20000). Thus, substrates can be LiF CuGaS $_2$ samples with varying substruc
109	ture including a super lattice phase super period.Growth mechanism of single-crystal
110	thin - film and nana scaleepitaxial films is a model for many heterogeneous and
111	topochemical processes.
112	However, the existing theory of crystallization cannot explain all of the results of a large a
113	mount of experimental work -
114	there is a clear discrepancy between the flow : The totality of new facts and their level of th
115	eoretical understanding of the experimentally observed facts with a unified position cannot
116	be considered. In existing theories of crystallization there is still no consensus on
117	what is themain factor in orienting epitaxial.

Generally accepted explanation is not considered as such a statement and that the main factor in orienting epitaxial is a single-crystal structure -

substrates. Focused on the growth of amorphous boundary layers prepared on the surface o
 f the substrate crystals [6,7], is proof that the basic structure of crystals does not

122 determine substrate orientation effects.

Oriented crystallization on the outside boundary of polycrystalline layers [8-123 9], as well as the growth of not only epitaxial films, but also highly perfect single 124 crystals through the amorphous boundary 125 126 layers [10] makes relate to the theory of crystallization based on the structure of single-127 crystal substrates seeds very carefully. Since many experimental works performed revealed aclear discrepancy between the facts and their level of theoretic 128 129 al understanding, it is not yet possible to formulate some general criteria for the formation o 130 f epitaxial films and single crystal. Therefore, conditions for the formation of singlecrystal films to day, as shown in previous chapter, 131 the are 132 established only experimentally.

133Since thebulk crystallattice CuGaS2ordered, in order to explain the134formation of super lattice phase should assume that it is the disordered phase. Disordering o135rdered - structuressome of the initial

136 of the atoms in it are defective, resulting superstructure should have a statistical average fr

equency.Regularities of similar phase transitions were first established for the phases of

138 the chemical group of compounds  $A^3B^3C_2^6[4]$ .

139 Superstructure phase  $CuGaS_2$  is

oriented (100) plane parallel to the faces [11] LiF. During epitaxial growth on LiF CuGaS<sub>2</sub>
one unit cell (UC) superstructure is

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142	mated with four cells of the substrate. Relative discrepancy mating crystal lattices in this
143	case is 2,9%.
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249 Fig.6.250



Fig. 7.

253 254 255 256 257	Captions to figuresofELECTRON DIFFRACTION AND ELECTRON MICROSCOPYSTUDY OF FILMS CuGaS <sub>2</sub> .
258 259	Fig.1. The intensity curves of amorphous CuGaS <sub>2</sub> .
260 261	Fig.2. Interference scattering function of electrons of amorphous $CuGaS_2$
262 263	Fig.3. The curve of radial distribution of atoms of $CuGaS_2$
264 265	Fig. 4. Electron diffraction single crystal mixture with polycrystalline $CuGaS_2$
266 267	Fig.5. Electron diffraction from single crystal super lattic phase CuGaS <sub>2</sub>
268 269	Fig.6. Electron diffraction pattern with Kikuchi lines from
270 271	single crystal CuGaS <sub>2</sub> high perfections.
272 273	Fig. 7. The microstructure of single-crystal beds of $CuGaS_2$ (H20000).